



Subspace-Based Near-Field Source Localization in Unknown Spatially Nonuniform Noise Environment

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Abstract—In this paper, we investigate the problem of estimating the directions-of-arrival (DOAs) and ranges of multiple narrow-band near-field sources in unknown spatially nonuniform noise (spatially inhomogeneous temporary white noise) environment, which is usually encountered in many practical applications of sensor array processing. A new subspace-based localization of near-field sources (SLONS) is proposed by exploiting the advantages of a symmetric uniform linear sensor array and using Toeplitzization of the array correlations. Firstly three Toeplitz correlation matrices are constructed by using the anti-diagonal elements of the array covariance matrix, where the nonuniform variances of additive noises are reduced to a uniform one, and then the location parameters (i.e., the DOAs and ranges) of near-field sources can be estimated by using the MUSIC-like method, while a new pair-matching scheme is developed to associate the estimated DOAs and ranges. Additionally, an alternating iterative scheme is considered to improve the estimation accuracy of the location parameters by utilizing the oblique projection operator, where the “saturation behavior” caused by finite number of snapshots is overcome effectively. Furthermore, the closed-form stochastic Cramér-Rao lower bound (CRB) is also derived explicitly for the near-field sources in the additive unknown nonuniform noises. Finally, the effectiveness of the proposed method and the theoretical analysis are substantiated through numerical examples.

Index Terms—Near-field, oblique projector, source localization, symmetric uniform linear array, unknown spatially nonuniform noises.

I. INTRODUCTION

LOCALIZATION of near-field sources is important in many practical application scenarios of sensor array processing

Manuscript received November 19, 2019; revised June 12, 2020 and July 14, 2020; accepted July 21, 2020. Date of publication August 5, 2020; date of current version September 3, 2020. The associate editor coordinating the review of this manuscript and approving it for publication was Dr. Abd-Krim Seghouane. This work was supported in part by the National Natural Science Foundation of China under Grants 61790563, 61671373, and 61627811, and in part by the National Key R&D Program of China under Grant 2017YFC0803905. This paper was presented in part at the IEEE 10th Sensor Array and Multichannel Signal Processing Workshop, Sheffield, U.K., July 8–11, 2018. (*Corresponding author: Jingmin Xin*)

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Digital Object Identifier 10.1109/TSP.2020.3013419

such as sonar, collision avoidance radar, electronic surveillance, seismology, speech enhancement, and biomedical imaging (e.g., [1]–[13] and references therein), where the near-field source is characterized by two independent location parameters (i.e., the directions-of-arrival (DOAs) and the ranges), and the correct pair-matching (i.e., association) of the estimated DOAs and ranges is required, while the far-field source is only parameterized by the DOA. In past decades, the localization of near-field sources was extensively studied (e.g., [14]–[34]). Among them, many maximum likelihood methods and subspace-based methods were proposed by using the second-order statistics (SOS) for spatially and temporally white Gaussian noises with the uniform variance or arbitrary noises with known statistics [14]–[32], while some subspace-based methods were presented by using the higher-order statistics (HOS) for spatially and temporally white or colored Gaussian noises [33], [34]. However, the homogeneous noise assumption is often invalid and the noise covariance matrix is usually unavailable in some applications (see, e.g., [35]–[42], [105]–[107]), which causes most subspace-based methods for the far-field sources (see, e.g., [43]–[47], [108]–[114]) to be severely degraded, and hence a variety of DOA estimation methods were developed for the far-field sources in the unknown colored noise [48]–[55] or unknown nonuniform noises [56]–[75]. Similarly in the presence of unknown nonuniform noises, most of the aforementioned SOS-based localization methods for the near-field sources will seriously deteriorate [14]–[32], while the HOS-based localization methods [33], [34] only apply to the non-Gaussian near-field sources. To the best of our knowledge, the localization of near-field sources in unknown spatially nonuniform noise environment has not been well studied in the literature of array processing. By using the multiple signal classification (MUSIC) [44], the two-dimensional (2D) MUSIC-whitened noise (MUSIC-WN) and the 2D MUSIC-smooth sparse arrays (MUSIC-SSA) were suggested in [76], where the noise-only array data or well-separated subarrays is required to whiten or eliminate the effect of additive noises, but a signal-free environment is not always available in practice and this requirement may restrict the practical applications (cf. [60], [54]). In addition, the far-field sources and the near-field sources usually coexist (see, e.g., [28] and references therein for details), and recently some HOS-based methods were proposed to localize the far-field and near-field sources in the presence of the additive white or colored Gaussian noises [77]–[84].

Therefore, in this paper, we deal with the localization of multiple near-field sources impinging on the symmetric uniform linear sensor array in the unknown spatially additive nonuniform noises (spatially inhomogeneous temporary white Gaussian noises), where the estimated DOAs should be associated with the corresponding estimated ranges of the same source. Based on our preliminary study [85], a new subspace-based localization of near-field sources (SLONS) is proposed by using the anti-diagonal correlation elements of the array covariance matrix, which can be interpreted as the received “signals” for a virtual ULA and contain only one noise variance or no noise variance. First by constructing two Toeplitz correlation matrices with the correlation elements along the major anti-diagonal or that along the 2nd upper diagonal off the major anti-diagonal of the array covariance matrix to eliminate the effects of unknown additive nonuniform noises, two auxiliary parameters related to the DOAs and ranges of near-field sources can be estimated separately from these two resulting matrices by using the MUSIC-like method [44]. As a result, we can obtain two sets of the estimated electric angles, which are functions of the location parameters (i.e., the DOAs and ranges). Then by forming another Toeplitz correlation matrix with the correlation elements along the 1st upper diagonal off the major anti-diagonal of the array covariance matrix, a pair-matching scheme is developed to associate two sets of the estimated electric angles by searching for the minimum value of the cost function, which is formed with the noise subspace of the resulting Toeplitz matrix, and consequently the estimated DOA and range of the same source can be obtained. Obviously unlike the pairing scheme presented in [23] for the near-field sources, which requires some constraints on the electric angles of sources, this new pair-matching scheme can applied to more general sources. Additionally an alternating iterative scheme is presented to improve the estimation accuracy of the location parameters by utilizing the oblique projection operator [86], where the “saturation behavior” (i.e., the estimated DOAs and ranges may have high elevated error floors, which do not decrease with the increasing signal-to-noise ratio (SNR)) caused by finite number of snapshots is overcome effectively. Furthermore, the stochastic Cramér-Rao lower bound (CRB) for the near-field sources in unknown nonuniform noises is derived explicitly. Finally, the effectiveness of the proposed method and the theoretical analysis are verified through numerical examples, and simulation results show that the proposed method achieves a remarkable estimation performance. Compared with the previous conference paper [85], the detailed theoretical derivation of the CRB is presented for the near-field sources in the unknown nonuniform noises, and the computational complexity of the proposed SLONS is also analyzed.

Glossary of Notation: The following notations are used through this paper. \mathbf{I}_m , $\mathbf{O}_{m \times n}$, $\mathbf{0}_{m \times 1}$, and $\delta_{n,t}$ stand for the $m \times m$ identity matrix, the $m \times n$ null matrix, the $m \times 1$ null vector, and the Kronecker delta, \mathbf{e}_l , $\bar{\mathbf{e}}_l$, and $\tilde{\mathbf{e}}_l$, stand for the $(M+1) \times 1$, $M \times 1$, and $(2M+1) \times 1$ unit vector with a unity element at the l th, l th, and $(M+1+l)$ th location and zeros elsewhere, and $E\{\cdot\}$, $\{\cdot\}^*$, $(\cdot)^T$, and $(\cdot)^H$ represent the statistical expectation, the complex conjugate, the transposition, and the Hermitian transposition, respectively. Additionally,

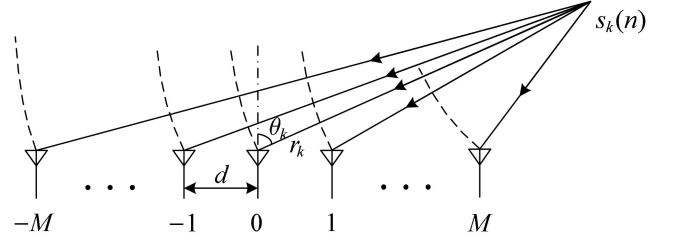


Fig. 1. The localization of near-field sources with a symmetric uniform linear sensor array.

$\text{diag}(\cdot)$, $\text{tr}\{\cdot\}$, $\text{rank}\{\cdot\}$, and $\mathcal{R}(\cdot)$ denote the diagonal matrix operator, the trace operator, the rank and the range space of the bracketed matrix, respectively, while \otimes , \odot , and \oplus signify the Kronecker product, the Hadamard-Schur product, and the direct sum operator. Furthermore $\text{vec}(\cdot)$ is a matrix operation to stack the columns of the bracketed matrix one under the other beginning with the leftmost column, $\mathcal{O}[\cdot]$ indicates the order of magnitude, $\text{span}\{\cdot\}$ denotes the columns space spanned by the column vectors of the bracketed matrix, and \hat{x} means the estimate of x .

II. DATA MODEL AND ASSUMPTIONS

As depicted in Fig. 1, we consider a symmetric uniform linear sensor array consisting of $2M+1$ sensors and K narrowband noncoherent signals $\{s_k(n)\}$ impinging on this array from the near-field sources with the unknown location parameters $\{\theta_k, r_k\}$, where θ_k is the DOA of the source signal $s_k(n)$ measured at the reference sensor relative to the normal of array, r_k is the corresponding range between the signal source and the reference sensor and given by $r_k \in (0.62(D^3/\lambda)^{1/2}, 2D^2/\lambda)$ [14], while D is the aperture of array given by $D = 2Md$ herein (cf. [4]), d is the sensor spacing, and λ is the wavelength. The received noisy signal $x_m(n)$ at the m th sensor can be approximated as

$$x_m(n) = \sum_{k=1}^K s_k(n) e^{j\tau_{mk}} + w_m(n) \quad (1)$$

for $m = -M, \dots, -1, 0, 1, \dots, M$, where $w_m(n)$ is the additive noise, and τ_{mk} is the phase delay due to the time delay between the reference sensor and the m th sensor for the signal $s_k(n)$ from the k th near-field source, which is given by (cf. [14])

$$\tau_{mk} = \frac{2\pi}{\lambda} \left(\sqrt{r_k^2 + (md)^2} - 2mdr_k \sin \theta_k - r_k \right). \quad (2)$$

Further, τ_{mk} in (2) can be approximated with the second-order Taylor expansion as [14], [23]

$$\tau_{mk} \approx m\psi_k + m^2\phi_k \quad (3)$$

where ψ_k and ϕ_k are called as the electric angles (include the DOAs and ranges) defined by

$$\psi_k \triangleq -\frac{2\pi d}{\lambda} \sin \theta_k \quad (4)$$

$$\phi_k \triangleq \frac{\pi d^2}{\lambda r_k} \cos^2 \theta_k. \quad (5)$$

Then the received data $\{x_m(n)\}$ can be rewritten compactly by using vector-matrix notation as

$$\mathbf{x}(n) = \mathbf{A}\mathbf{s}(n) + \mathbf{w}(n) \quad (6)$$

where $\mathbf{x}(n)$, $\mathbf{s}(n)$, and $\mathbf{w}(n)$ are the vectors of the received data, the incident signals and the additive noises given by $\mathbf{x}(n) \triangleq [x_{-M}(n), \dots, x_{-1}(n), x_0(n), x_1(n), \dots, x_M(n)]^T$, $\mathbf{s}(n) \triangleq [s_1(n), s_2(n), \dots, s_K(n)]^T$, and $\mathbf{w}(n) \triangleq [w_{-M}(n), \dots, w_{-1}(n), w_0(n), w_1(n), \dots, w_M(n)]^T$, while \mathbf{A} is the array response matrix defined by $\mathbf{A} \triangleq [\mathbf{a}(\theta_1, r_1), \mathbf{a}(\theta_2, r_2), \dots, \mathbf{a}(\theta_K, r_K)]$, and $\mathbf{a}(\theta_k, r_k)$ is the array steering vectors which can be expressed as $\mathbf{a}(\theta_k, r_k) \triangleq [e^{-jM\psi_k} e^{jM^2\phi_k}, \dots, e^{-j\psi_k} e^{j\phi_k}, 1, e^{j\psi_k} e^{j\phi_k}, \dots, e^{jM\psi_k} e^{jM^2\phi_k}]^T$.

Here we make the basic assumptions as follows.

- A1) The array is calibrated and the array response matrix \mathbf{A} has full rank and is unambiguous.
- A2) Without loss of generality, the incident signals $\{s_k(n)\}$ from the K near-field sources are temporally complex white Gaussian random processes with zero-mean and the variance given by $E\{s_k(n)s_k^*(t)\} = r_{s_k} \delta_{n,t}$ and $E\{s_k(n)s_k(t)\} = 0, \forall n, t$.
- A3) For the simplicity of theoretical performance analysis, the additive noises $\{w_i(n)\}$ are temporally and spatially complex Gaussian random processes with zero-mean, and the covariance matrix of unknown additive noises is given by

$$\mathbf{Q} \triangleq E\{\mathbf{w}(n)\mathbf{w}^H(n)\} \\ = \text{diag}(\sigma_{-M}^2, \dots, \sigma_{-1}^2, \sigma_0^2, \sigma_1^2, \dots, \sigma_M^2) \quad (7)$$

where σ_m^2 is the noise power at the m th sensor. Moreover, the covariance matrices $E\{\mathbf{w}(n)\mathbf{w}^T(t)\} = \mathbf{O}_{(2M+1) \times (2M+1)}, \forall n, t$. Additionally the additive noises $\{w_m(n)\}$ are independent to the incident signals $\{s_k(n)\}$, i.e., $E\{\mathbf{s}(n)\mathbf{w}^H(n)\} = E\{\mathbf{s}(n)\mathbf{w}^T(n)\} = \mathbf{O}_{K \times (2M+1)}$.

- A4) The number of near-field sources K is known, and K satisfies the relation $K < M$ (cf. Remark C for details).
- A5) The sensor spacing d satisfies the relation $d \leq \lambda/4$ for avoiding the estimation ambiguity.

In the following, we concentrate on the estimation of location parameters $\{\theta_k\}_{k=1}^K$ and $\{r_k\}_{k=1}^K$ of multiple near-field sources from the finite noisy array data $\{\mathbf{x}(n)\}_{n=1}^N$, where the basic idea is firstly to transfer the nonuniform noise into an uniform one and then to estimate the location parameters from the reformed data model under the virtual uniform noise environment.

Remark A: Various localization techniques have been developed in many application fields (see [6], [7] and references therein), where the time-of-arrival (TOA), the time-difference-of-arrival (TDOA), the received signal strength (RSS), and the DOA are commonly used measurements for source localization (e.g., [87]–[90]). Basically the TOA, TDOA, and RSS provide the distance information between the source and the sensors, while the DOA gives the source bearing relative to the sensors. Perhaps the TDOA is the most widely used technique for passive

localization [90], it measures the difference in time between the source signal received by a pair of spatially separated sensors (i.e., two receivers), where the sensors are required to be synchronized to a common time reference. By multiplying the TDOA with a known propagation speed, the range difference between the source and two sensors can be obtained. As a result, the source location is given by the intersection of at least two hyperbolas, where a hyperbola is defined by the range difference between the source and two sensors. Obviously it will be rather difficult and complicated to localize multiple sources by using the TDOA technique, though the TDOA localization was studied for one source in far-field and near-field [89]. In this paper, we consider the localization of multiple near-field source by estimating their DOAs and ranges. \square

III. NEW SUBSPACE-BASED LOCALIZATION OF NEAR-FIELD SOURCES—SLONS

Here we develop a method called SLONS for estimating the DOAs and ranges of multiple near-field sources by exploiting the special configuration of a symmetric uniform linear array. The proposed SLONS includes three stages: 1) estimate the location parameters (i.e., DOAs and ranges) of the near-fields sources, 2) associate the estimated DOAs and ranges, and 3) performance improvement against finite snapshots.

A. Direction-of-Arrival Estimation of Near-Field Sources

From (6), we obtain the array covariance matrix \mathbf{R} as

$$\mathbf{R} \triangleq E\{\mathbf{x}(n)\mathbf{x}^H(n)\} = \mathbf{A}\mathbf{R}_s\mathbf{A}^H + \mathbf{Q} \quad (8)$$

where $\mathbf{R}_s \triangleq E\{\mathbf{s}(n)\mathbf{s}^H(n)\} = \text{diag}(r_{s_1}, r_{s_2}, \dots, r_{s_K})$ with $r_{s_k} \triangleq E\{s_k(n)s_k^*(n)\}$, and the anti-diagonal correlation element of \mathbf{R} is given by (cf. [23])

$$R_i(p) \triangleq E\{x_p(n)x_{-p-i}^*(n)\} \\ = \mathbf{r}_{is}^T \mathbf{b}_i(p) + \sigma_{0.5i}^2 \delta_{p,-p-i} \quad (9)$$

for $i = -2M, \dots, -1, 0, 1, \dots, 2M$ and $p = -M + i_2, \dots, -1, 0, 1, \dots, M - i_1$, where $i_1 = 0.5(|i| + i)$, $i_2 = 0.5(|i| - i)$, $\mathbf{b}_i(p) \triangleq [e^{j2p\gamma_{i1}}, e^{j2p\gamma_{i2}}, \dots, e^{j2p\gamma_{iK}}]^T$, $\mathbf{r}_{is} \triangleq [r_{s_1} e^{j\gamma_{i1}}, r_{s_2} e^{j\gamma_{i2}}, \dots, r_{s_K} e^{j\gamma_{iK}}]^T$, $\gamma_{ik} = \psi_k - i\phi_k$, and $\sigma_{0.5i}^2 = 0$ when i is odd. Herein we easily have

$$\psi_k = \gamma_{0k} \quad (10)$$

$$\phi_k = \frac{1}{2}(\gamma_{0k} - \gamma_{2k}). \quad (11)$$

Obviously $\{R_i(p)\}$ can be interpreted as the received “signals” for a virtual array of $2M + 1 - |i|$ sensors illuminated by K “signals” $\{r_{s_k} e^{j\gamma_{ik}}\}$, and $\{R_i(p)\}$ differ only by a phase factor γ_{ik} . Hence the auxiliary parameter γ_{ik} can be estimated by using the phase delays of $\{R_i(p)\}$.

By setting $i = 0$ (i.e., $i_1 = i_2 = 0$), we have $p = -M, -M + 1, \dots, -1, 0, 1, \dots, M - 1, M$. Therefore, by using the $2M + 1$ correlations $\{R_0(p)\}_{p=-M}^M$ along the major cross-diagonal of \mathbf{R} (i.e., $i = 0$), we can form an $(M + 1) \times (M + 1)$ Toeplitz

matrix \mathbf{T}_0 as

$$\mathbf{T}_0 = \begin{bmatrix} R_0(0), & R_0(-1), & \cdots, & R_0(-M) \\ R_0(1), & R_0(0), & \cdots, & R_0(-M+1) \\ \vdots & \vdots & \ddots & \vdots \\ R_0(M), & R_0(M-1), & \cdots, & R_0(0) \end{bmatrix} \quad (12)$$

where its (l, t) th element can be represented as

$$(\mathbf{T}_0)_{l,t} = R_0(l-t) \quad (13)$$

for $l, t = 1, 2, \dots, M+1$, in which the t th column \mathbf{c}_{0t} of \mathbf{T}_0 is given by

$$\mathbf{c}_{0t} \triangleq [(\mathbf{T}_0)_{1,t}, (\mathbf{T}_0)_{2,t}, \dots, (\mathbf{T}_0)_{M+1,t}]^T. \quad (14)$$

By substituting (9) into (14), after some simple manipulations, we have

$$\begin{aligned} \mathbf{c}_{0t} &= [R_0(1-t), R_0(2-t), \dots, R_0(M+1-t)]^T \\ &= [\mathbf{b}_0(1-t), \mathbf{b}_0(2-t), \dots, \mathbf{b}_0(M+1-t)]^T \mathbf{r}_{0s} \\ &\quad + [\mathbf{0}_{(t-1) \times 1}^T, \sigma_0^2, \mathbf{0}_{(M+1-t) \times 1}^T]^T \\ &= \mathbf{A}_0 \mathbf{D}_0^{1-t} \mathbf{r}_{0s} + \sigma_0^2 \mathbf{e}_t \end{aligned} \quad (15)$$

where the Kronecker delta $\delta_{l-t, -l+t} = 1$ for $l-t = -l+t$ (i.e., $l = t$) is used, and $\mathbf{D}_0 = \text{diag}(e^{j2\gamma_{01}}, e^{j2\gamma_{02}}, \dots, e^{j2\gamma_{0K}})$, while $\mathbf{A}_0 \triangleq [\mathbf{b}_0(0), \mathbf{b}_0(1), \dots, \mathbf{b}_0(M)]^T$ with the dimension $(M+1) \times K$. Then from (13)–(15), the Toeplitz correlation matrix \mathbf{T}_0 in (12) can be rewritten as

$$\begin{aligned} \mathbf{T}_0 &= [\mathbf{c}_{01}, \mathbf{c}_{02}, \dots, \mathbf{c}_{0,M+1}] \\ &= \mathbf{A}_0 [\mathbf{R}_s \mathbf{b}_0(0), \mathbf{R}_s \mathbf{b}_0^*(1), \dots, \mathbf{R}_s \mathbf{b}_0^*(M)] + \sigma_0^2 \mathbf{I}_{M+1} \\ &= \mathbf{A}_0 \mathbf{R}_s \mathbf{A}_0^H + \sigma_0^2 \mathbf{I}_{M+1}. \end{aligned} \quad (16)$$

Clearly the nonuniform noise matrix \mathbf{Q} in (8) is transformed to a uniform one $\sigma_0^2 \mathbf{I}_{M+1}$ in (16), and we can see that the matrix \mathbf{A}_0 is a Vandermonde matrix with full rank, while the eigenvalue decomposition (EVD) of the matrix \mathbf{T}_0 in (16) is given by

$$\mathbf{T}_0 = \mathbf{U}_{0s} \mathbf{A}_{0s} \mathbf{U}_{0s}^H + \mathbf{U}_{0n} \mathbf{A}_{0n} \mathbf{U}_{0n}^H \quad (17)$$

where \mathbf{A}_{0s} and \mathbf{A}_{0n} are the diagonal matrices consisting of the K largest and the $M+1-K$ smallest eigenvalues respectively, while the $(M+1) \times K$ signal subspace matrix \mathbf{U}_{0s} and the $(M+1) \times (M+1-K)$ noise subspace matrix \mathbf{U}_{0n} consist of the corresponding eigenvectors, respectively. The columns of the Vandermonde matrix \mathbf{A}_0 and the columns of the signal subspace matrix \mathbf{U}_{0s} span the same subspace, i.e., $\text{span}\{\mathbf{A}_0\} = \text{span}\{\mathbf{U}_{0s}\}$. As a result, we have

$$\mathbf{U}_{0n}^H \mathbf{A}_0 = \mathbf{O}_{(M+1) \times K}. \quad (18)$$

Therefore by using the relation (18), we can easily get the estimates of the auxiliary parameters $\{\gamma_{0k}\}_{k=1}^K$ (i.e., the DOAs $\{\theta_k\}_{k=1}^K$) with the conventional MUSIC method [44]. Furthermore, when the number of snapshots is finite, the auxiliary parameters $\{\gamma_{0k}\}_{k=1}^K$ can be estimated by minimizing the following MUSIC-like cost function (cf. [44])

$$f_0(\gamma_{0k}) = \mathbf{a}_0^H(\gamma_{0k}) \hat{\mathbf{U}}_{0n} \hat{\mathbf{U}}_{0n}^H \mathbf{a}_0(\gamma_{0k}) \quad (19)$$

where $\mathbf{a}_0(\gamma_{0k})$ denotes the k th column of \mathbf{A}_0 given by $\mathbf{a}_0(\gamma_{0k}) = [1, e^{j2\gamma_{0k}}, \dots, e^{j2M\gamma_{0k}}]^T$. Then from (4) and (10), we can find that the DOAs of the near-field sources can be estimate from the auxiliary parameter $\{\gamma_{0k}\}_{k=1}^K$ (see details in Section III-C).

B. Range Estimation of Near-Field Sources

Similarly by setting $i = 2$ (i.e., $i_1 = 2, i_2 = 0$, and $p = -M, -M+1, \dots, -1, 0, 1, \dots, M-3, M-2$) and by using the $2M-1$ correlations $\{R_2(p)\}_{p=-M}^{M-2}$ along the 2nd upper diagonal off the major cross-diagonal of \mathbf{R} , we can form another $M \times M$ Toeplitz matrix \mathbf{T}_2 as

$$\begin{aligned} \mathbf{T}_2 &= \begin{bmatrix} R_2(-1), & R_2(-2), & \cdots, & R_2(-M) \\ R_2(0), & R_2(-1), & \cdots, & R_2(-M+1) \\ \vdots & \vdots & \ddots & \vdots \\ R_2(M-2), & R_2(M-3), & \cdots, & R_2(-1) \end{bmatrix} \\ &= \mathbf{A}_2 \mathbf{R}_s \mathbf{A}_2^H + \sigma_M^2 \mathbf{I}_M \end{aligned} \quad (20)$$

where the (l, t) th element of \mathbf{T}_2 can be represented as

$$(\mathbf{T}_2)_{l,t} = R_2(l-t-1) \quad (21)$$

for $l, t = 1, 2, \dots, M$, and $\mathbf{A}_2 \triangleq [\mathbf{b}_2(0), \mathbf{b}_2(1), \dots, \mathbf{b}_2(M-1)]^T$. Obviously the nonuniform noise matrix \mathbf{Q} in (8) is transformed to a uniform one $\sigma_1^2 \mathbf{I}_M$ in (20), and the matrix \mathbf{A}_2 is a Vandermonde matrix, which has full rank. Hence the EVD of the matrix \mathbf{T}_2 in (20) is given by

$$\mathbf{T}_2 = \mathbf{U}_{2s} \mathbf{A}_{2s} \mathbf{U}_{2s}^H + \mathbf{U}_{2n} \mathbf{A}_{2n} \mathbf{U}_{2n}^H \quad (22)$$

where \mathbf{A}_{2s} and \mathbf{A}_{2n} are the diagonal matrices consisting of the K largest and the $M-K$ smallest eigenvalues respectively, while the $M \times K$ signal subspace matrix \mathbf{U}_{2s} and the $M \times (M-K)$ noise subspace matrix \mathbf{U}_{2n} consists of the corresponding eigenvectors, respectively. Then we easily have

$$\mathbf{U}_{2n}^H \mathbf{A}_2 = \mathbf{O}_{M \times K}. \quad (23)$$

When the number of snapshots is finite, the auxiliary parameters $\{\gamma_{2k}\}_{k=1}^K$ can be estimated by minimizing the following MUSIC-like cost function (cf. [44])

$$f_2(\gamma_{2k}) = \mathbf{a}_2^H(\gamma_{2k}) \hat{\mathbf{U}}_{2n} \hat{\mathbf{U}}_{2n}^H \mathbf{a}_2(\gamma_{2k}) \quad (24)$$

where $\mathbf{a}_2(\gamma_{2k})$ denotes the k th column of \mathbf{A}_2 given by $\mathbf{a}_2(\gamma_{2k}) = [1, e^{j2\gamma_{2k}}, \dots, e^{j2M\gamma_{2k}}]^T$.

Although two sets of the auxiliary parameters $\{\hat{\gamma}_{0k}\}_{k=1}^K$ and $\{\hat{\gamma}_{2k}\}_{k=1}^K$ are estimated separately with (19) and (24), there is no clear one-to-one relationship between each element of these two estimation sets. From (10) and (11), we easily get K estimates $\{\hat{\psi}_1, \hat{\psi}_2, \dots, \hat{\psi}_K\}$ of the electric angle ψ_k (a function of DOA θ_k as shown in (4)) and K^2 temporary estimates $\{\hat{\phi}_{11}, \hat{\phi}_{12}, \dots, \hat{\phi}_{1K}, \dots, \hat{\phi}_{KK}\}$ of the electric angle ϕ_k (which is a function of DOA θ_k and range r_k) by

$$\hat{\psi}_k = \hat{\gamma}_{0k}, \quad \hat{\phi}_{kk'} = 0.5(\hat{\gamma}_{0k} - \hat{\gamma}_{2k'}) \quad (25)$$

for $k = 1, 2, \dots, K$ and $k' = 1, 2, \dots, K$. In order to estimate and associate the DOA and range of the same near-field source from these two separate sets of estimates $\{\hat{\psi}_k\}_{k=1}^K$ and

$\{\hat{\phi}_{kk'}\}_{k,k'=1}^K$, an additional procedure of pair-matching is required.

C. Parameter Pairing and Estimation of DOA and Range

Although two sets of the estimated electric angles $\hat{\psi}_k$ and $\hat{\phi}_{kk'}$ are obtained in (25) for $k = 1, 2, \dots, K$ and $k' = 1, 2, \dots, K$, they should be associated with a corresponding source to estimate the location parameters (i.e., the DOA and range) of the same near-field source. Now we consider a new pair-matching scheme for the DOA and range estimation.

Herein by setting $i = 1$ (i.e., $i_1 = 1, i_2 = 0, p = -M, -M + 1, \dots, -1, 0, 1, \dots, M - 2, M - 1$) and by employing the $2M - 1$ correlations $\{R_1(p)\}_{p=-M+1}^{M-1}$ along the 1st upper diagonal off the major cross-diagonal of \mathbf{R} in (9), we can form a new $M \times M$ Toeplitz matrix \mathbf{T}_1 as

$$\mathbf{T}_1 = \begin{bmatrix} R_1(0), & R_1(-1), & \dots, & R_1(-M+1) \\ R_1(1), & R_1(0), & \dots, & R_1(-M+2) \\ \vdots & \vdots & \ddots & \vdots \\ R_1(M-1), & R_1(M-2), & \dots, & R_1(0) \end{bmatrix} \\ = \mathbf{A}_1 \bar{\mathbf{R}}_s \mathbf{A}_1^H \quad (26)$$

where the (l, t) th element of \mathbf{T}_1 can be represented as

$$(\mathbf{T}_1)_{l,t} = R_1(l - t) \quad (27)$$

for $l, t = 1, 2, \dots, M$, $\mathbf{A}_1 \triangleq [\mathbf{b}_1(0), \mathbf{b}_1(1), \dots, \mathbf{b}_1(M-1)]^T$, and $\bar{\mathbf{R}}_s = \text{diag}(r_{s_1} e^{j\gamma_{i1}}, r_{s_2} e^{j\gamma_{i2}}, \dots, r_{s_K} e^{j\gamma_{iK}})$. Then the EVD of the matrix \mathbf{T}_1 is given by

$$\mathbf{T}_1 = \mathbf{U}_{1s} \mathbf{A}_{1s} \mathbf{U}_{1s}^H + \mathbf{U}_{1n} \mathbf{A}_{1n} \mathbf{U}_{1n}^H \quad (28)$$

where \mathbf{A}_{1s} and \mathbf{A}_{1n} are the diagonal matrices consisting of the K largest eigenvalues (i.e., nonzero) and the $M - K$ smallest eigenvalues (i.e., zeros) respectively, while \mathbf{U}_{1s} and \mathbf{U}_{1n} are the $M \times K$ signal subspace and the $M \times (M - K)$ noise subspace, respectively. When the number of snapshots is finite, we can estimate $\{\gamma_{1k}\}_{k=1}^K$ by minimizing the following cost function

$$f_1(\gamma_{1k}) = \mathbf{a}_{1k}^H(\gamma_{1k}) \hat{\mathbf{U}}_{1n} \hat{\mathbf{U}}_{1n}^H \mathbf{a}_{1k}(\gamma_{1k}) \quad (29)$$

where $\mathbf{a}_{1k}(\gamma_{1k})$ denotes the k th column of \mathbf{A}_1 given by $\mathbf{a}_{1k}(\gamma_{1k}) = [1, e^{j2\gamma_{1k}}, \dots, e^{j2M\gamma_{1k}}]^T$. Hence from (25) and (29), the estimated electric phase angles $\hat{\psi}_{k_o}$ and $\hat{\phi}_{kk'}$ can be associated with the indices (k_o, k'_o) determined by

$$(k_o, k'_o) = \arg \min_{(k, k')} f_1(\hat{\psi}_k, \hat{\phi}_{kk'}) \quad (30)$$

for $k, k' = 1, 2, \dots, K$.

Finally by using these paired estimates $\hat{\psi}_{k_o}$ and $\hat{\phi}_{k_o, k'_o}$ (i.e., $\hat{\psi}_k$ and $\hat{\phi}_k$), from (4) and (5), the DOA θ_k and range r_k of the incident signal $s_k(n)$ are obtained as

$$\hat{\theta}_k = \arcsin(\alpha_1 \hat{\psi}_k) \quad (31)$$

$$\hat{r}_k = \frac{\alpha_2}{\hat{\phi}_k} \cos^2(\hat{\theta}_k) = \frac{\alpha_2}{\hat{\phi}_k} \cos^2(\arcsin(\alpha_1 \hat{\psi}_k)) \quad (32)$$

where $\alpha_1 = -\lambda/(2\pi d)$, and $\alpha_2 = \pi d^2/\lambda$. As a result, the location parameters (i.e., $\{\hat{\theta}_k, \hat{r}_k\}_{k=1}^K$) of the same near-field source can be estimated.

D. Alternating Iterative Scheme Against Finite Snapshots

When the number of snapshots is not sufficiently large enough, the nonzero residual cross-correlations between the incident signals will cause the estimated matrices $\hat{\mathbf{R}}_s$ and $\hat{\mathbf{R}}$ (and hence the Toeplitz matrices $\hat{\mathbf{T}}_0, \hat{\mathbf{T}}_1$, and $\hat{\mathbf{T}}_2$) be inaccurate. Consequently the estimated DOA $\hat{\theta}_k$ and range \hat{r}_k may have high elevated error floors, which do not decrease with the increasing SNR (i.e., “saturation behavior”) (cf. [91]–[93]). Since the range space of each incident signal is nonoverlapping and not orthogonal to that of the others, we can use the oblique projection operator [86] to isolate one signal from the others and present an alternating iterative scheme to improve the performance with finite snapshots of array data.

By defining $\mathbf{a}_k \triangleq \mathbf{a}(\theta_k, r_k)$ and denoting the reduced signal vector without $s_k(n)$ and the corresponding reduced array response matrix without column \mathbf{a}_k as $\bar{\mathbf{s}}_k(n)$ and $\bar{\mathbf{A}}_k$, where $\mathcal{R}(\mathbf{A}) = \mathcal{R}(\mathbf{a}_k) \oplus \mathcal{R}(\bar{\mathbf{A}}_k)$, we can obtain

$$\hat{\mathbf{R}}_s = \begin{bmatrix} \hat{r}_{s_k}, & \hat{\rho}^T \\ \hat{\rho}^*, & \hat{\mathbf{R}}_k \end{bmatrix} \quad (33)$$

$$\hat{\mathbf{R}} = \hat{r}_{s_k} \mathbf{a}_k \mathbf{a}_k^H + \mathbf{a}_k \hat{\rho}^T \bar{\mathbf{A}}_k^H + \bar{\mathbf{A}}_k \hat{\rho}^* \mathbf{a}_k^H + \bar{\mathbf{A}}_k \hat{\mathbf{R}}_k \bar{\mathbf{A}}_k^H + \hat{\mathbf{R}}_e \quad (34)$$

where $\hat{\mathbf{R}}_e = \mathbf{A} \hat{\mathbf{R}}_{sw} + \hat{\mathbf{R}}_{sw}^H \mathbf{A}^H + \hat{\mathbf{Q}}$, $\hat{\rho} = (1/N) \sum_{n=1}^N s_k(n) \cdot \bar{\mathbf{s}}_k(n)$, $\hat{\mathbf{R}}_k = (1/N) \sum_{n=1}^N \bar{\mathbf{s}}_k(n) \bar{\mathbf{s}}_k^H(n)$, while $\hat{\mathbf{R}}_{sw} = (1/N) \cdot \sum_{n=1}^N \mathbf{s}(n) \mathbf{s}^H(n)$, and the oblique projector $\mathbf{E}_{\bar{\mathbf{A}}_k|\mathbf{a}_k}$ which projects onto the space $\mathcal{R}(\bar{\mathbf{A}}_k)$ along a direction parallel to the space $\mathcal{R}(\mathbf{a}_k)$ is given by (cf. [86], [91])

$$\mathbf{E}_{\bar{\mathbf{A}}_k|\mathbf{a}_k} \triangleq \bar{\mathbf{A}}_k (\bar{\mathbf{A}}_k^H \mathbf{\Pi}_{\mathbf{a}_k}^\perp \bar{\mathbf{A}}_k)^{-1} \bar{\mathbf{A}}_k^H \mathbf{\Pi}_{\mathbf{a}_k}^\perp \quad (35)$$

where $\mathbf{\Pi}_{\mathbf{a}_k}^\perp \triangleq \mathbf{I}_{2M+1} - \mathbf{a}_k (\mathbf{a}_k^H \mathbf{a}_k)^{-1} \mathbf{a}_k^H$, $\mathbf{E}_{\bar{\mathbf{A}}_k|\mathbf{a}_k} \bar{\mathbf{A}}_k = \bar{\mathbf{A}}_k$, and $\mathbf{E}_{\bar{\mathbf{A}}_k|\mathbf{a}_k} \mathbf{a}_k = \mathbf{O}_{(2M+1) \times 1}$. Because the oblique projector with true parameters in (35) is not available, by denoting the estimated electrical phase angles as $\{\hat{\psi}_k^{(t)}\}_{k=1}^K$ and $\{\hat{\phi}_k^{(t)}\}_{k=1}^K$, we can get an estimated oblique projector as $\hat{\mathbf{E}}_{\bar{\mathbf{A}}_k|\mathbf{a}_k}^{(t)}$. Then, from (34) and (35), we get the projected correlation matrix $\hat{\mathbf{R}}_k^{(t)}$ as

$$\hat{\mathbf{R}}_k^{(t)} = (\mathbf{I}_{2M+1} - \hat{\mathbf{E}}_{\bar{\mathbf{A}}_k|\mathbf{a}_k}^{(t)}) \hat{\mathbf{R}} (\mathbf{I}_{2M+1} - \hat{\mathbf{E}}_{\bar{\mathbf{A}}_k|\mathbf{a}_k}^{(t)})^H \\ \approx \hat{r}_{s_k} \hat{\mathbf{a}}_k \hat{\mathbf{a}}_k^H. \quad (36)$$

Obviously $\hat{\mathbf{R}}_k^{(t)}$ only contains the information of the signal $s_k(n)$. Thus, by estimating the electrical phase angles as discussed in Sections III.A and updating the index as $t = t + 1$, we can obtain the renewed estimates $\{\hat{\psi}_k^{(t+1)}\}_{k=1}^K$ and $\{\hat{\phi}_k^{(t+1)}\}_{k=1}^K$. This procedure should be repeated several times until the difference between two consecutive iterations becomes smaller than

a threshold, i.e.,

$$\sum_{k=1}^K \left| \hat{\psi}_k^{(t+1)} - \hat{\psi}_k^{(t)} \right| \leq \varepsilon \quad (37)$$

where ε is an arbitrary and positive small constant (e.g., $\varepsilon = 10^{-4}$), then finally we can obtain the refined estimates $\hat{\theta}_k$ and \hat{r}_k of the source signal $s_k(n)$ from these $\hat{\psi}_k^{(t+1)}$ and $\hat{\phi}_k^{(t+1)}$.

Therefore when the finite array data are available, the proposed SLONS method can be implemented as follows:

- 1) Estimate the array covariance matrix as $\hat{\mathbf{R}} = (1/N) \cdot \sum_{n=1}^N \mathbf{x}(n)\mathbf{x}^H(n)$.
..... $10(2M+1)^2 N$ flops
- 2) Estimate $\{\gamma_{0k}\}_{k=1}^K$ by finding the phases of the K zeros of the polynomial $p_0(z)$ closest to the unit circle in the z -plane, where $p_0(z) \triangleq \mathbf{p}_0^H(z) \hat{\mathbf{U}}_{0n} \hat{\mathbf{U}}_{0n}^H \mathbf{p}_0(z)$, $\mathbf{p}_0(z) \triangleq [1, z^{-1}, \dots, z^{-M}]^T$, and $z \triangleq e^{j2\gamma_0}$.
..... $\mathcal{O}[(2M+1)^2 + (M+1)^3]$ flops
- 3) Estimate $\{\gamma_{2k}\}_{k=1}^K$ by finding the phases of the K zeros of the polynomial $p_2(z)$ closest to the unit circle in the z -plane, where $p_2(z) \triangleq \mathbf{p}_2^H(z) \hat{\mathbf{U}}_{2n} \hat{\mathbf{U}}_{2n}^H \mathbf{p}_2(z)$, $\mathbf{p}_2(z) \triangleq [1, z^{-1}, \dots, z^{-(M-1)}]^T$, and $z \triangleq e^{j2\gamma_2}$.
..... $\mathcal{O}[(2M-1)^2 + M^3]$ flops
- 4) Estimate and associate the electrical phase angles with (25) and (30), and denote them as $\{\hat{\psi}_k^{(t)}\}_{k=1}^K$ and $\{\hat{\phi}_k^{(t)}\}_{k=1}^K$, where $t = 0$.
..... $16M^2 + 8M^2(M-K) + \mathcal{O}[M^3]$ flops
- 5) For $k = 1, 2, \dots, K$, by calculating the oblique projector $\mathbf{E}_{\mathcal{A}_k|\mathcal{A}_k}^{(t)}$ with (35) from $\{\hat{\psi}_k^{(t)}\}_{k=1}^K$ and $\{\hat{\phi}_k^{(t)}\}_{k=1}^K$, and the projected matrix $\hat{\mathbf{R}}_k^{(t)}$ with (36), estimate the renewed electrical phase angles $\{\hat{\psi}_k^{(t+1)}\}_{k=1}^K$ and $\{\hat{\phi}_k^{(t+1)}\}_{k=1}^K$ with (19) and (29)-(30).
..... $16(2M+1)^3 + (2M+1)^2(24K-23) + 8(2M+1)(2(K-1)^2+1) + \mathcal{O}[(2M+1)^2 + (M+1)^3]$ flops
- 6) If the condition in (37) is not satisfied, repeat Step 5 by setting $t = t+1$; otherwise by reexpressing the estimates $\{\hat{\psi}_k^{(t+1)}\}_{k=1}^K$ and $\{\hat{\phi}_k^{(t+1)}\}_{k=1}^K$ as $\{\hat{\psi}_k\}_{k=1}^K$ and $\{\hat{\phi}_k\}_{k=1}^K$, estimate the DOAs $\{\hat{\theta}_k\}_{k=1}^K$ and ranges $\{\hat{r}_k\}_{k=1}^K$ with (31) and (32).

The computational complexity of each step is roughly indicated in terms of the number of MATLAB flops, where a flop is defined as a floating-point addition or multiplication operation,¹ and hence the computational complexity of the SLONS is approximately $10(2M+1)^2 N + 16(2M+1)^3 N_i$ if $2M+1 \gg K$ flops, which occurs often in applications of array processing, where N_i denotes the times of iteration ($N_i \geq 1$).

Remark B: The quantitative comparisons of computational complexities between the SLONS and some existing localization methods for the near-field sources in the additive uniform noise

¹Such as the computational complexities of $\mathbf{X} \pm \mathbf{Y}$ and $\mathbf{X}\mathbf{Y}$ (\mathbf{X} and \mathbf{Y} are $L \times M$ and $M \times N$ complex matrices, respectively) are $2LM$ and $8LMN$, respectively. See more details about the calculation of the computational complexity in [47] and references therein.

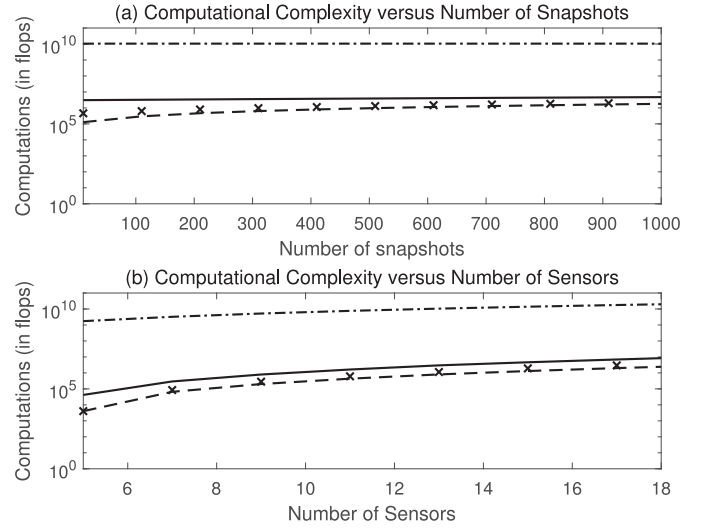


Fig. 2. Comparison of computational complexities in MATLAB flops versus (a) the number of snapshots ($2M+1=11$) and (b) the number of sensors ($N=200$) (dot-dashed line: 2D-MUSIC; "x": WLPM; solid line: the SLONS with 10 iterations; and dashed line: the SLONS w/o iteration).

such as the modified 2D-MUSIC [15] and the weighted linear prediction method (WLPM) [23] are shown in Fig. 2. The 2D-MUSIC is based on the spherical wavefront model and involves 2D spectrum peak searching, hence its computational complexity is about $10(2M+1)^2 N + (8(2M+1)^2 + 8(2M+1))(180/\Delta\theta)((2D^2/\lambda - 0.62(D^3/\lambda)^{1/2})/\Delta r) + \mathcal{O}[(2M+1)^3]$ flops, where $\Delta\theta$ and Δr are the angular and range grid spacings (i.e., $\Delta\theta = 0.002^\circ$ and $\Delta r = 0.002\lambda$), and the precise peak searching necessitates fine grid spacing, but it is a rather time-consuming task with heavy computational load. On the contrary, the WLPM is based on the Fresnel approximation of the spherical wavefront model, and its computational complexity is reduced roughly to $10(2M+1)^2 N + \mathcal{O}[(2M+1-K)^3]$ flops. From Fig. 2, we can see that the proposed SLONS is computationally more efficient than the 2D-MUSIC and similar to the WLPM. \square

Remark C: The $2M+1$, $2M$ and $2M-1$ correlations along three different anti-diagonals of $\hat{\mathbf{R}}$ are used to construct three Toeplitz matrices with the dimensions $(M+1) \times (M+1)$, $M \times M$, and $M \times M$, and hence the maximum resolvable signals of the SLONS is $K < M$. This is roughly the same as the WLPM [23] and other subspace-based near-field localization methods (e.g., [24], [92], [93]), where the geometric symmetry of the uniform linear sensor array is exploited to facilitate the parameter estimation (cf. [94]). \square

IV. STOCHASTIC CRB FOR NEAR-FIELD SOURCES IN UNKNOWN SPATIALLY NONUNIFORM NOISE

The CRB provides a lower bound on the variance of any unbiased estimator (see, e.g., [95]). Although the expressions of the stochastic and deterministic CRBs were studied for the far-field sources under the nonuniform noise assumption [57], [100], [101], and some expressions of the stochastic and deterministic CRBs were derived for the near-field sources

in the spatially uniform noise environment [16], [23], [96]–[99], to the best of our knowledge, the CRB for the near-field source localization in the spatially nonuniform noises is still unavailable so far. Therefore, we derive the closed-form stochastic CRB for the estimated DOAs and ranges of the near-field sources in the presence of spatially nonuniform noises herein.

Under the data model assumptions, the Fisher information matrix (FIM) for a parameter vector α is given by (e.g., [95], [100]–[102])

$$(\text{FIM})_{k,l} = N \text{tr} \left\{ \frac{\partial \mathbf{R}}{\partial \alpha_k} \mathbf{R}^{-1} \frac{\partial \mathbf{R}}{\partial \alpha_l} \mathbf{R}^{-1} \right\} \quad (38)$$

for $k, l = 1, 2, \dots, 3K + 2M + 1$, where $\alpha \triangleq [\zeta^T, \rho^T, \sigma^T]^T$, $\zeta \triangleq [\theta_1, \theta_2, \dots, \theta_K, r_1, r_2, \dots, r_K]^T$, $\rho \triangleq [r_{s1}, r_{s2}, \dots, r_{sK}]^T$, and $\sigma \triangleq [\sigma_1^2, \sigma_2^2, \dots, \sigma_{2M+1}^2]^T$. To focus only on the desired ζ -block of the $\text{CRB} = \text{FIM}^{-1}$, which is denoted by $\text{CRB}(\zeta)$, we can obtain a closed-form expression of the CRB for the estimated DOAs and ranges of the near-field sources as the following theorem.

Theorem 1: The $2K \times 2K$ stochastic CRB for the location parameters of the near-field sources under unknown spatially nonuniform noise environment is given by

$$\text{CRB}(\zeta) = \frac{1}{N} (\mathbf{F} - \mathbf{H} \mathbf{K}^{-1} \mathbf{H}^T)^{-1} \quad (39)$$

where

$$\mathbf{F} = 2\text{Re} \left\{ (\tilde{\mathbf{D}}^H \Pi_{\tilde{\mathbf{A}}}^{\perp} \tilde{\mathbf{D}}) \odot (\mathbf{J} \otimes (\mathbf{R}_s \tilde{\mathbf{A}}^H \tilde{\mathbf{R}}^{-1} \tilde{\mathbf{A}} \mathbf{R}_s)^T) \right\} \quad (40)$$

$$\mathbf{H} = 2\text{Re} \left\{ (\tilde{\mathbf{D}}^H \Pi_{\tilde{\mathbf{A}}}^{\perp}) \odot (\mathbf{1} \otimes (\tilde{\mathbf{R}}^{-1} \tilde{\mathbf{A}} \mathbf{R}_s)^T) \right\} \quad (41)$$

$$\mathbf{K} = 2\text{Re} \left\{ \Pi_{\tilde{\mathbf{A}}}^{\perp} \odot \tilde{\mathbf{R}}^{-T} \right\} - \Pi_{\tilde{\mathbf{A}}}^{\perp} \odot (\Pi_{\tilde{\mathbf{A}}}^{\perp})^T \quad (42)$$

while $\tilde{\mathbf{R}} \triangleq \mathbf{Q}^{-1/2} \mathbf{R} \mathbf{Q}^{-1/2}$, $\tilde{\mathbf{A}} \triangleq \mathbf{Q}^{-1/2} \mathbf{A}$, $\tilde{\mathbf{D}} = \mathbf{Q}^{-1/2} \mathbf{D}$, $\Pi_{\tilde{\mathbf{A}}}^{\perp} \triangleq \mathbf{I}_{2M+1} - \tilde{\mathbf{A}}(\tilde{\mathbf{A}}^H \tilde{\mathbf{A}})^{-1} \tilde{\mathbf{A}}^H$, $\mathbf{1} \triangleq [1, 1]^T$, $\mathbf{J} \triangleq \mathbf{1} \mathbf{1}^T$, and $\mathbf{D} \triangleq [d_{\theta_1}, d_{\theta_2}, \dots, d_{\theta_K}, d_{r_1}, d_{r_2}, \dots, d_{r_K}]$, in which $d_{\theta_k} \triangleq \partial \mathbf{a}(\theta_k, r_k) / \partial \theta_k$ and $d_{r_k} \triangleq \partial \mathbf{a}(\theta_k, r_k) / \partial r_k$.

Proof: See Appendix. ■

V. NUMERICAL EXAMPLES

Here we evaluate the effectiveness of the proposed SLONS through numerical examples. The uniform linear sensor array consists of $2M + 1 = 11$ sensors with element spacing $d = \lambda/4$, while the complex incident signals $\{s_k(n)\}$ are generated by

$$s_k(n) = s_k^r(n) + j s_k^i(n) \quad (43)$$

where the real part $s_k^r(n)$ and the imaginary part $s_k^i(n)$ are two different random processes given by Matlab Random Number Generator function “randn”. The SNR is defined as

$$\text{SNR} = 10 \log_{10} \frac{1}{2M+1} \sum_{m=-M}^M \frac{\sigma_m^2}{r_s} \quad (44)$$

where the incident signals have the same power, i.e., $r_{s1} = \dots = r_{sK} = r_s$. The 2D-MUSIC [15] and the WLPM [23] are carried

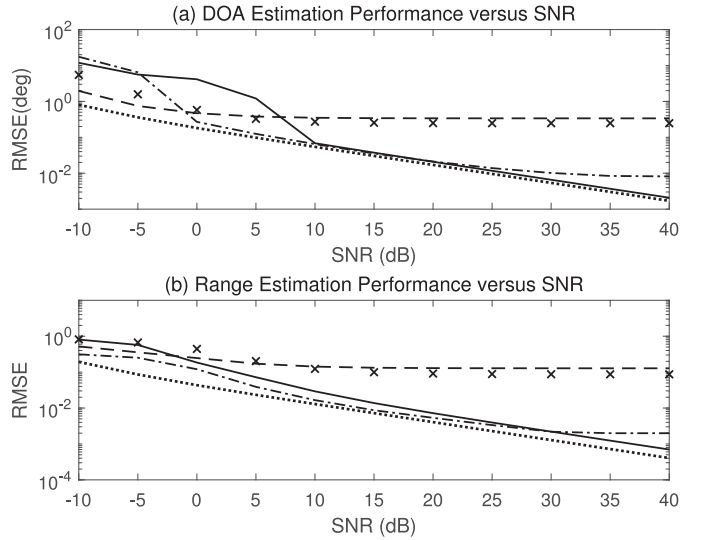


Fig. 3. RMSE of the estimated DOAs and ranges versus the SNR (“x”: WLPM; dot-dashed line: 2D-MUSIC; dashed line: the SLONS w/o iteration; solid line: the SLONS; dotted line: CRB) in Example 1.

out, and the stochastic CRB derived in (39) is calculated for performance comparison. The angular and range grid spacings $\Delta\theta$ and Δr are fixed at $\Delta\theta = 0.09^\circ$ and $\Delta r = 0.09\lambda$ for the 2D-MUSIC, the iteration threshold in (37) is set at $\varepsilon = 10^{-6}$. The results are all based on 1000 independent trials (i.e., $\bar{N} = 1,000$).

Example 1—Performance versus SNR: There are three near-field sources with equal power located at $(-11^\circ, 2.9\lambda)$, $(21^\circ, 3.0\lambda)$, and $(47^\circ, 3.3\lambda)$. The number of snapshots is $N = 500$, and the SNR is varied from -10 dB to 40 dB, where the covariance matrix of additive noises is given by

$$\mathbf{Q} = \sigma_n \text{diag}(2, 2, 0.7, 1, 1.5, 5, 3, 3.6, 4, 2, 3). \quad (45)$$

The averaged empirical root mean-squared-errors (RMSEs) of the estimated location parameters in terms of the SNR are plotted in Figs. 3(a) and 3(b). The 2D-MUSIC [15] utilizes all $(2M + 1)^2$ correlations of $\hat{\mathbf{R}}$, and it is rather time-consuming but its performance is restricted by the angular/range grid spacing, while the WLPM [23] uses the $2M + 1$, $2M$ and $2M$ correlations along three anti-diagonals of $\hat{\mathbf{R}}$ with a rather low computational cost, but it is rather complicated to determine the optimal weighting matrix. Moreover, the structure properties of the array covariance matrices of the incident signals and the additive noise are required in the 2D-MUSIC and the WLPM, but the structure properties are only valid for large number of snapshots. When the number of snapshots is not sufficiently large enough, the erroneous estimated correlations will cause the linear prediction (LP) model to be invalid, and correspondingly the WLPM suffers severely “saturation behavior” regardless of the SNR as shown in Fig. 3, where the estimated DOAs and ranges have high elevate error floors and will not decrease with the increasing SNR. The 2D-MUSIC has relatively good performance at low and moderate SNRs, but the estimation errors of the DOAs and ranges no longer decrease at high SNR

due to the limited searching grids of the 2D spectrum. Although the saturation behavior of the 2D-MUSIC could be alleviated by a refined searching grid, the amount of computational load will increase dramatically. In this paper, since the range space of each incident signal is nonoverlapping and not orthogonal to that of the others, the oblique projection operator (see, e.g., [86]) is used to isolate one signal from the others, and an alternating iterative scheme is presented to improve the performance with finite snapshots of array data, where the location parameters (i.e., two electric angles) of one source are estimated in one iteration, and the mutual influence between the multiple near-field sources are eliminated. As shown in Figs. 3(a) and 3(b), we can find that the SLONS has a good tradeoff between the computational complexity and the estimation performance compared with the 2D-MUSIC [15] and is superior to the WLPM [23], where its empirical RMSEs of the estimated DOAs and ranges are very close to the CRBs at moderate and high SNRs, and the “saturation behavior” is overcome effectively. Unfortunately, the SLONS performs poorly at relatively low SNR because only some (but not all) correlations are used compared with the 2D-MUSIC.

Example 2—Performance versus Number of Snapshots: The simulation conditions are similar to those in Example 1, except that the SNR is fixed at 20 dB and the number of snapshots varies from 100 to 10,000.

Figs. 4(a) and 4(b) show the averaged empirical RMSEs of the estimated DOAs and ranges with respect to the number of snapshots. Since the WLPM uses the similar correlations along three anti-diagonals of the array covariance matrix, its performance is similar to the SLONS without iteration, and although the RMSEs of the estimated DOAs and ranges decrease with the increasing number of snapshots, these RMSEs cannot reach the corresponding CRBs due to “saturation behavior”. Additionally as described in Example 1, owing to the limited searching grips of the 2D spectrum, the RMSEs of the 2D-MUSIC deviate from the CRBs even for larger number of snapshots. As shown in Figs. 4(a) and 4(b), when the number of snapshots is not large, the estimates initial values of the electric angles used for the iteration may have a large deviation, and the iteration may not converge, hence the SLONS will have relatively larger estimation error. However when the number of snapshots is larger, the SLONS with iteration outperforms the WLPM for both the estimation of DOAs and ranges, where the “saturation behavior” is alleviated.

Example 3—Convergence Behavior: The simulation conditions are similar to those in above examples, and the averaged convergence behaviors of the proposed alternating iterative scheme are plotted in Fig. 5 for different SNR and number of snapshots.

From Fig. 5, we can see that the alternating iterative scheme converges within roughly 20 times of iterations and the “saturation behavior” can be resolved for small number of snapshots, while the convergence can be obtained in a few iterations and the performance is improved remarkably for sufficiently large number of snapshots. However, the convergence analysis of the alternating iterative scheme is rather complicated, and the

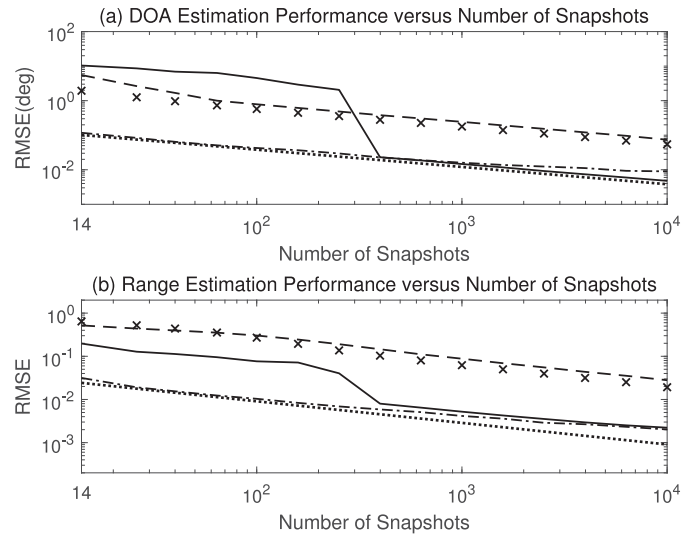


Fig. 4. RMSE of the estimated DOAs and ranges versus the number of snapshots (“x”: WLPM; dot-dashed line: 2D-MUSIC; dashed line: the SLONS w/o iteration; solid line: the SLONS; dotted line: CRB) in Example 2.

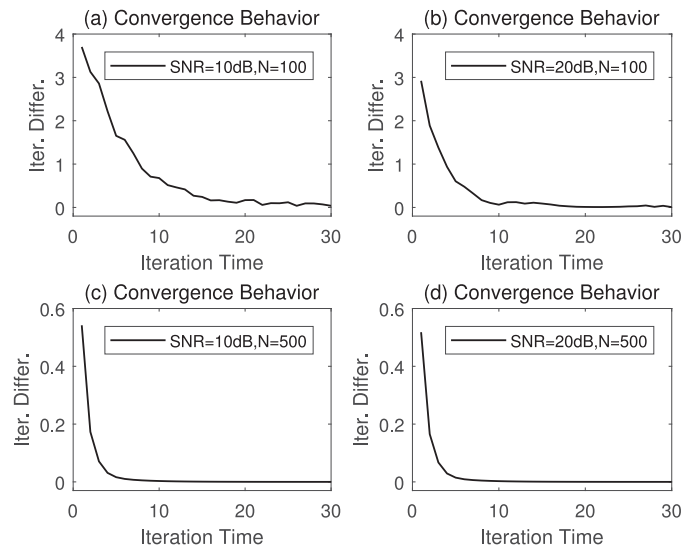


Fig. 5. Averaged convergence behaviors of the alternating iterative scheme for several SNRs and numbers of snapshots in Example 3.

convergence depends on the initialization of $\{\hat{\psi}_k^{(0)}\}_{k=1}^K$ and $\{\hat{\phi}_k^{(0)}\}_{k=1}^K$.

Example 4—Performance versus Angular Separation: In this example, two near-field sources are located at $(11^\circ, 2.9\lambda)$ and $(11^\circ + \Delta\theta, 2.9\lambda)$, where $\Delta\theta$ is varied from 0.5° to 19.5° with $\Delta\theta = 1^\circ$, and the SNR is fixed at 15 dB, and the number of snapshots is set at $N = 500$.

Fig. 6 displays the averaged empirical RMSEs of the estimated DOAs and ranges against the angular separation. We can find that the proposed SLONS generally outperforms the WLPM and the 2D-MUSIC for the closely-spaced sources, and the RMSEs of the location parameters do not decrease monotonically with

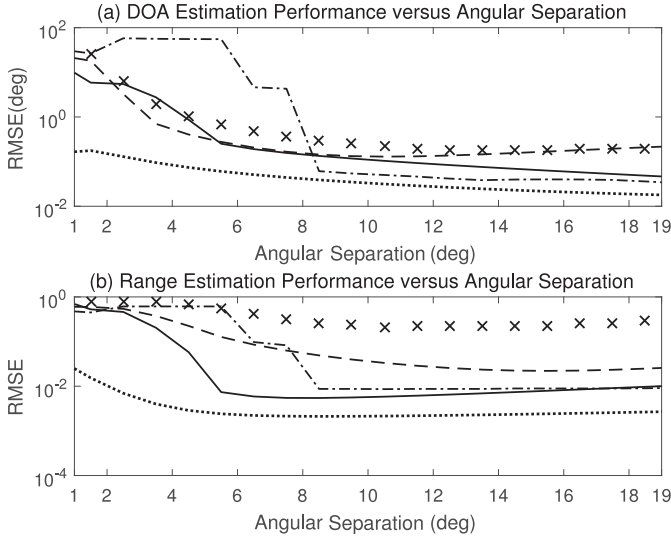


Fig. 6. RMSEs of (a) the estimated DOAs versus the angular separation and (b) the estimated ranges versus the angular separation for Example 4 (“x”: WLPM; dot-dashed line: 2D-MUSIC; dashed line: the SLONS w/o iteration; solid line: the SLONS; dotted line: CRB) (SNR = 15dB, $N = 500$, $(11^\circ, 2.9\lambda)$, and $(11^\circ + \Delta\theta, 2.9\lambda)$).

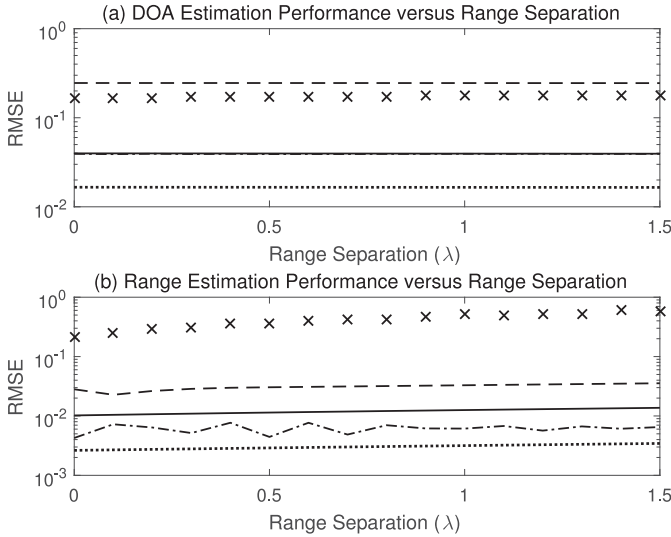


Fig. 7. RMSEs of (a) the estimated DOAs versus the range separation and (b) the estimated ranges versus the range separation for Example 4 (“x”: WLPM; dot-dashed line: 2D-MUSIC; dashed line: the SLONS w/o iteration; solid line: the SLONS; dotted line: CRB) (SNR = 15dB, $N = 500$, $(11^\circ, 2.9\lambda)$, and $(11^\circ + \Delta\theta, 2.9\lambda + \Delta\lambda)$).

the increasing angular separation. It is also worth mentioning that the proposed SLONS is computationally more efficient than the 2D-MUSIC, though its performance of the DOA estimation is worse than the 2D-MUSIC for some large angular separations.

Example 5—Performance versus Range Separation: The simulation conditions are similar to those in Example 4, while two near-field sources are located at $(11^\circ, 2.9\lambda)$ and $(11^\circ, 2.9\lambda + \Delta\lambda)$, where $\Delta\lambda$ is varied from 0λ to 1.5λ with $\Delta\lambda = 0.1\lambda$.

Fig. 7 demonstrates the averaged empirical RMSEs of the estimated location parameters in terms of the range separation. When the source range increases, the electric angle ϕ_k in (5) will become very small, and the estimation performance of ϕ_k will be an ill-posed problem [23]. Hence the range estimation performance will become worse from (32), while the DOA estimation performance is almost constant, where the estimation performance is mainly dominated by the number of snapshots and SNR. Obviously the estimation performance of the proposed SLONS is effectively improved by using the iteration, and the SLONS is superior to the WLPM for both DOA and range estimation. Additionally the SLONS has similar DOA estimation performance but poor range estimation performance compared with the 2D-MUSIC in this empirical scenario.

It should be noted that the empirical RMSEs of the DOA and range estimates do not decrease monotonically as the range separation increases. Furthermore the SLONS can be extended to the multiple far-field sources in presence of unknown nonuniform noises, where the source location is characterized only by the DOA.

VI. CONCLUSION

A new subspace-based method called SLONS was proposed for localization of multiple near-field sources impinging on a symmetric uniform linear sensor array in the unknown spatially additive nonuniform noises. A new pari-matching scheme was developed, and an oblique projector based alternating iterative scheme was presented to combat the “saturation behavior” due to finite number of snapshots. Further the closed-form stochastic CRB was derived for the localization of near-field sources in unknown nonuniform noises explicitly. The simulation results showed that the proposed method has remarkable estimation performance.

APPENDIX PROOF OF THEOREM 1

The derivation of the CRB will use the following readily-checked facts (see, e.g., [104])

$$\text{tr}\{\mathbf{X}\mathbf{Y}\} = \text{vec}^H(\mathbf{X}^H)\text{vec}(\mathbf{Y}) \quad (\text{A1})$$

$$\text{vec}\{\mathbf{X}\mathbf{Y}\mathbf{Z}\} = (\mathbf{Z}^T \otimes \mathbf{X})\text{vec}(\mathbf{Y}) \quad (\text{A2})$$

which hold for any conformable matrices \mathbf{X} , \mathbf{Y} , and \mathbf{Z} . By using (A1) and (A2), the matrix FIM associated with the data model in (6) can be expressed as (cf. [101], [102])

$$\text{FIM} = N \left(\frac{d\mathbf{y}}{d\boldsymbol{\alpha}^T} \right)^H (\mathbf{R}^{-T} \otimes \mathbf{R}^{-1}) \left(\frac{d\mathbf{y}}{d\boldsymbol{\alpha}^T} \right) \quad (\text{A3})$$

where $\mathbf{y} = \text{vec}\{\mathbf{R}\} = (\mathbf{A}^* \otimes \mathbf{A})\text{vec}(\mathbf{R}_s) + \text{vec}(\mathbf{Q})$. By defining the following two matrices as

$$\mathbf{G} \triangleq (\mathbf{R}^{-T/2} \otimes \mathbf{R}^{-1/2}) \frac{d\mathbf{y}}{d\boldsymbol{\zeta}^T} \quad (\text{A4})$$

$$\boldsymbol{\Delta} \triangleq (\mathbf{R}^{-T/2} \otimes \mathbf{R}^{-1/2}) \left[\frac{d\mathbf{y}}{d\boldsymbol{\rho}^T}, \frac{d\mathbf{y}}{d\boldsymbol{\sigma}^T} \right] \quad (\text{A5})$$

from (A3), we can obtain

$$\frac{1}{N} \text{FIM} = \begin{bmatrix} \mathbf{G}^H \\ \Delta^H \end{bmatrix} [\mathbf{G}, \Delta] \quad (\text{A6})$$

and hence the CRB for the parameter ζ consisting of the DOAs and ranges of near-field sources under unknown nonuniform noise environment is given by (cf. [101], [102])

$$\text{CRB}(\zeta) = \frac{1}{N} \text{FIM}^{-1} = \frac{1}{N} (\mathbf{G}^H \Pi_{\Delta}^{\perp} \mathbf{G})^{-1} \quad (\text{A7})$$

where $\Pi_{\Delta}^{\perp} = \mathbf{I} - \Pi_{\Delta}$, and $\Pi_{\Delta} \triangleq \Delta(\Delta^H \Delta)^{-1} \Delta^H$, while Π_{Δ} and Π_{Δ}^{\perp} stand for the orthogonal projection matrices onto the range space of Δ and its orthogonal complement, respectively, and the inversion of partitioned matrices [103] is used implicitly. Further by partitioning the matrix Δ in (A5) into two submatrices as

$$\Delta = [\mathbf{V}, \mathbf{U}] \quad (\text{A8})$$

where

$$\mathbf{V} \triangleq (\mathbf{R}^{-T/2} \otimes \mathbf{R}^{-1/2}) \frac{d\mathbf{y}}{d\rho^T} \quad (\text{A9})$$

$$\mathbf{U} \triangleq (\mathbf{R}^{-T/2} \otimes \mathbf{R}^{-1/2}) \frac{d\mathbf{y}}{d\sigma^T}. \quad (\text{A10})$$

Since the matrices Δ and $[\mathbf{V}, \Pi_{\mathbf{V}}^{\perp} \mathbf{U}]$ have the same range space, Π_{Δ}^{\perp} in (A7) is easily given by (cf. [101], [102])

$$\Pi_{\Delta}^{\perp} = \Pi_{\mathbf{V}}^{\perp} - \Pi_{\mathbf{V}}^{\perp} \mathbf{U} (\mathbf{U}^H \Pi_{\mathbf{V}}^{\perp} \mathbf{U})^{-1} \mathbf{U}^H \Pi_{\mathbf{V}}^{\perp}. \quad (\text{A11})$$

Now by reexpressing the parameter ζ in (38) as $\zeta = [\theta^T, \mathbf{r}^T]^T$, where $\theta \triangleq [\theta_1, \theta_2, \dots, \theta_K]^T$, and $\mathbf{r} \triangleq [r_1, r_2, \dots, r_K]^T$, the matrix \mathbf{G} in (A4) can be rewritten as

$$\mathbf{G} = (\mathbf{R}^{-T/2} \otimes \mathbf{R}^{-1/2}) \begin{bmatrix} \frac{d\mathbf{y}}{d\theta^T} & \frac{d\mathbf{y}}{d\mathbf{r}^T} \end{bmatrix} = [\mathbf{G}_{\theta}, \mathbf{G}_r] \quad (\text{A12})$$

where

$$\mathbf{G}_{\theta} \triangleq (\mathbf{R}^{-T/2} \otimes \mathbf{R}^{-1/2}) \frac{d\mathbf{y}}{d\theta^T} \quad (\text{A13})$$

$$\mathbf{G}_r \triangleq (\mathbf{R}^{-T/2} \otimes \mathbf{R}^{-1/2}) \frac{d\mathbf{y}}{d\mathbf{r}^T}. \quad (\text{A14})$$

Then by substituting (A8)-(A14) into (A7), after some manipulations, we can obtain

$$\begin{aligned} \text{CRB}(\zeta) &= \frac{1}{N} (\mathbf{F} - \mathbf{H} \mathbf{K}^{-1} \mathbf{H}^T)^{-1} \\ &= \frac{1}{N} \begin{bmatrix} \mathbf{F}_{\theta\theta} - \mathbf{H}_{\theta} \mathbf{K}^{-1} \mathbf{H}_{\theta}^H & \mathbf{F}_{\theta r} - \mathbf{H}_{\theta} \mathbf{K}^{-1} \mathbf{H}_r^H \\ \mathbf{F}_{\theta r}^H - \mathbf{H}_r \mathbf{K}^{-1} \mathbf{H}_{\theta}^H & \mathbf{F}_{rr} - \mathbf{H}_r \mathbf{K}^{-1} \mathbf{H}_r^H \end{bmatrix}^{-1} \end{aligned} \quad (\text{A15})$$

where

$$\mathbf{F} \triangleq \mathbf{G}^H \Pi_{\mathbf{V}}^{\perp} \mathbf{G} \quad (\text{A16})$$

$$\mathbf{H} \triangleq \mathbf{G}^H \Pi_{\mathbf{V}}^{\perp} \mathbf{U} \quad (\text{A17})$$

$$\mathbf{K} \triangleq \mathbf{U}^H \Pi_{\mathbf{V}}^{\perp} \mathbf{U} \quad (\text{A18})$$

while

$$\mathbf{F}_{\theta\theta} \triangleq \mathbf{G}_{\theta}^H \Pi_{\mathbf{V}}^{\perp} \mathbf{G}_{\theta} \quad (\text{A19})$$

$$\mathbf{F}_{\theta r} \triangleq \mathbf{G}_{\theta}^H \Pi_{\mathbf{V}}^{\perp} \mathbf{G}_r \quad (\text{A20})$$

$$\mathbf{F}_{rr} \triangleq \mathbf{G}_r^H \Pi_{\mathbf{V}}^{\perp} \mathbf{G}_r \quad (\text{A21})$$

$$\mathbf{H}_{\theta} \triangleq \mathbf{G}_{\theta}^H \Pi_{\mathbf{V}}^{\perp} \mathbf{U} \quad (\text{A22})$$

$$\mathbf{H}_r \triangleq \mathbf{G}_r^H \Pi_{\mathbf{V}}^{\perp} \mathbf{U} \quad (\text{A23})$$

where the property $\mathbf{F}_{r\theta} = \mathbf{F}_{\theta r}^H$ is used in (A15).

Further we need to calculate the derivatives of \mathbf{y} with respect to $\{\theta_k, r_k\}$. First, we consider $d\mathbf{y}/d\theta^T$ and $d\mathbf{y}/d\mathbf{r}^T$. Let \mathbf{p}_k denote the k th column of signal covariance matrix \mathbf{R}_s , i.e. $\mathbf{R}_s = [\mathbf{p}_1, \mathbf{p}_2, \dots, \mathbf{p}_K]$. Thus, we have

$$\frac{d\mathbf{R}}{d\theta_k} = d_{\theta_k} \mathbf{p}_k^H \mathbf{A}^H + \mathbf{A} \mathbf{p}_k d_{\theta_k} \quad (\text{A24})$$

$$\frac{d\mathbf{R}}{dr_k} = d_{r_k} \mathbf{p}_k^H \mathbf{A}^H + \mathbf{A} \mathbf{p}_k d_{r_k} \quad (\text{A25})$$

Therefore, the k th column of \mathbf{G} (i.e. \mathbf{G}_{θ} and \mathbf{G}_r) is given by

$$\mathbf{g}_{\theta_k} = \text{vec} \left(\mathbf{R}^{-1/2} \frac{d\mathbf{R}}{d\theta_k} \mathbf{R}^{-1/2} \right) \triangleq \text{vec}(\mathbf{Z}_{\theta_k} + \mathbf{Z}_{\theta_k}^H) \quad (\text{A26})$$

$$\mathbf{g}_{r_k} = \text{vec} \left(\mathbf{R}^{-1/2} \frac{d\mathbf{R}}{dr_k} \mathbf{R}^{-1/2} \right) \triangleq \text{vec}(\mathbf{Z}_{r_k} + \mathbf{Z}_{r_k}^H) \quad (\text{A27})$$

where

$$\mathbf{Z}_{\theta_k} = \mathbf{R}^{-1/2} \mathbf{A} \mathbf{p}_k d_{\theta_k}^H \mathbf{R}^{-1/2} \quad (\text{A28})$$

$$\mathbf{Z}_{r_k} = \mathbf{R}^{-1/2} \mathbf{A} \mathbf{p}_k d_{r_k}^H \mathbf{R}^{-1/2} \quad (\text{A29})$$

Then we consider $d\mathbf{y}/d\rho^T$. The derivation of this part is similar to the far-field case (e.g., [57]). Thus by using the following properties

$$\Pi_{\mathbf{V}}^{\perp} = \Pi_{(\mathbf{R}^{-1/2} \mathbf{A})^* \otimes (\mathbf{R}^{-1/2} \mathbf{A})}^{\perp} \quad (\text{A30})$$

$$\Pi_{(\mathbf{X} \otimes \mathbf{Y})} = \mathbf{I} \otimes \Pi_{\mathbf{Y}}^{\perp} + \Pi_{\mathbf{X}}^{\perp} \otimes \mathbf{I} - \Pi_{\mathbf{X}}^{\perp} \otimes \Pi_{\mathbf{Y}}^{\perp} \quad (\text{A31})$$

From (A26)-(A31), we obtain

$$\Pi_{\mathbf{V}}^{\perp} \mathbf{g}_{\theta_k} = \text{vec} (\Pi_{\mathbf{R}^{-1/2} \mathbf{A}}^{\perp} \mathbf{Z}_{\theta_k}^H + \mathbf{Z}_{\theta_k} \Pi_{\mathbf{R}^{-1/2} \mathbf{A}}^{\perp}) \quad (\text{A32})$$

$$\Pi_{\mathbf{V}}^{\perp} \mathbf{g}_{r_k} = \text{vec} (\Pi_{\mathbf{R}^{-1/2} \mathbf{A}}^{\perp} \mathbf{Z}_{r_k}^H + \mathbf{Z}_{r_k} \Pi_{\mathbf{R}^{-1/2} \mathbf{A}}^{\perp}) \quad (\text{A33})$$

where the following facts are used

$$\Pi_{\mathbf{R}^{-1/2} \mathbf{A}}^{\perp} \mathbf{Z}_{\theta_k} = \mathbf{O}_{(2M+1) \times (2M+1)} \quad (\text{A34})$$

$$\Pi_{\mathbf{R}^{-1/2} \mathbf{A}}^{\perp} \mathbf{Z}_{r_k} = \mathbf{O}_{(2M+1) \times (2M+1)} \quad (\text{A35})$$

Finally, from (A19)-(A21) along with (A1), (A26), (A27), (A32), and (A33), the (i, k) th element of \mathbf{F} can be obtained as

$$(\mathbf{F}_{\theta\theta})_{i,k} = 2\text{Re}\{\text{tr}(\mathbf{Z}_{\theta_i} \Pi_{\mathbf{R}^{-1/2} \mathbf{A}}^{\perp} \mathbf{Z}_{\theta_k}^H)\} \quad (\text{A36})$$

$$(\mathbf{F}_{\theta r})_{i,k} = 2\text{Re}\{\text{tr}(\mathbf{Z}_{\theta_i} \Pi_{\mathbf{R}^{-1/2} \mathbf{A}}^{\perp} \mathbf{Z}_{r_k}^H)\} \quad (\text{A37})$$

$$(\mathbf{F}_{rr})_{i,k} = 2\text{Re}\{\text{tr}(\mathbf{Z}_{r_i} \Pi_{\mathbf{R}^{-1/2} \mathbf{A}}^{\perp} \mathbf{Z}_{r_k}^H)\} \quad (\text{A38})$$

Furthermore, we can easily have

$$\mathbf{R}^{-1/2} \Pi_{\mathbf{R}^{-1/2} \mathbf{A}}^{\perp} \mathbf{R}^{-1/2} = \mathbf{Q}^{-1/2} \Pi_{\mathbf{A}}^{\perp} \mathbf{Q}^{-1/2} \triangleq \tilde{\Pi}_{\mathbf{A}}^{\perp} \quad (\text{A39})$$

By inserting (A28) and (A29) into (A36)–(A38) and then using (A39) yields

$$(\mathbf{F}_{\theta\theta})_{i,k} = 2\text{Re}\{(\mathbf{d}_{\theta_i} \tilde{\Pi}_{\tilde{\mathbf{A}}}^\perp \mathbf{d}_{\theta_k})(\mathbf{p}_k^H \mathbf{A}^H \mathbf{R}^{-1} \mathbf{A} \mathbf{p}_i)\} \quad (\text{A40})$$

$$(\mathbf{F}_{\theta r})_{i,k} = 2\text{Re}\{(\mathbf{d}_{\theta_i} \tilde{\Pi}_{\tilde{\mathbf{A}}}^\perp \mathbf{d}_{r_k})(\mathbf{p}_k^H \mathbf{A}^H \mathbf{R}^{-1} \mathbf{A} \mathbf{p}_i)\} \quad (\text{A41})$$

$$(\mathbf{F}_{rr})_{i,k} = 2\text{Re}\{(\mathbf{d}_{r_i} \tilde{\Pi}_{\tilde{\mathbf{A}}}^\perp \mathbf{d}_{r_k})(\mathbf{p}_k^H \mathbf{A}^H \mathbf{R}^{-1} \mathbf{A} \mathbf{p}_i)\} \quad (\text{A42})$$

Next, we consider the matrices \mathbf{H}_θ and \mathbf{H}_r . From (A22)–(A23) along with (A1), (A2), (A10), (A28), (A29), (A32), (A33), and (A39), the (i,k) th elements of \mathbf{H}_θ and \mathbf{H}_r are given by

$$(\mathbf{H}_\theta)_{i,k} = 2\text{Re}\{\tilde{\mathbf{e}}_m^T \mathbf{R}^{-1} \mathbf{A} \mathbf{p}_i \mathbf{d}_{\theta_k}^H \tilde{\Pi}_{\tilde{\mathbf{A}}}^\perp \tilde{\mathbf{e}}_m\} \quad (\text{A43})$$

$$(\mathbf{H}_r)_{i,k} = 2\text{Re}\{\tilde{\mathbf{e}}_m^T \mathbf{R}^{-1} \mathbf{A} \mathbf{p}_i \mathbf{d}_{r_k}^H \tilde{\Pi}_{\tilde{\mathbf{A}}}^\perp \tilde{\mathbf{e}}_m\} \quad (\text{A44})$$

where the fact $d\mathbf{R}/d\sigma_m = d\mathbf{Q}/d\sigma_m = \tilde{\mathbf{e}}_m \tilde{\mathbf{e}}_m^T$ is used in (A43) and (A44).

At last, we consider the matrix \mathbf{K} in (A18). By using (A1), (A2), (A10), (A30), (A31), and (A39), we can get the (i,k) th element of \mathbf{K} as follows

$$(\mathbf{K})_{i,k} = 2\text{Re}\{\tilde{\mathbf{e}}_m^T \tilde{\Pi}_{\tilde{\mathbf{A}}}^\perp \tilde{\mathbf{e}}_m \tilde{\mathbf{e}}_m^T \mathbf{R}^{-1} \tilde{\mathbf{e}}_m\} - (\tilde{\mathbf{e}}_m^T \tilde{\Pi}_{\tilde{\mathbf{A}}}^\perp \tilde{\mathbf{e}}_m)^2 \quad (\text{A45})$$

Then, by performing some direct and tedious manipulations, we easily get

$$\mathbf{F}_{\theta\theta} = 2\text{Re}\left\{(\tilde{\mathbf{D}}_\theta^H \Pi_{\tilde{\mathbf{A}}}^\perp \tilde{\mathbf{D}}_\theta) \odot (\mathbf{R}_s \tilde{\mathbf{A}}^H \tilde{\mathbf{R}}^{-1} \tilde{\mathbf{A}} \mathbf{R}_s)^T\right\} \quad (\text{A46})$$

$$\mathbf{F}_{\theta r} = 2\text{Re}\left\{(\tilde{\mathbf{D}}_\theta^H \Pi_{\tilde{\mathbf{A}}}^\perp \tilde{\mathbf{D}}_r) \odot (\mathbf{R}_s \tilde{\mathbf{A}}^H \tilde{\mathbf{R}}^{-1} \tilde{\mathbf{A}} \mathbf{R}_s)^T\right\} \quad (\text{A47})$$

$$\mathbf{F}_{r\theta} = 2\text{Re}\left\{(\tilde{\mathbf{D}}_r^H \Pi_{\tilde{\mathbf{A}}}^\perp \tilde{\mathbf{D}}_\theta) \odot (\mathbf{R}_s \tilde{\mathbf{A}}^H \tilde{\mathbf{R}}^{-1} \tilde{\mathbf{A}} \mathbf{R}_s)^T\right\} \quad (\text{A48})$$

$$\mathbf{F}_{rr} = 2\text{Re}\left\{(\tilde{\mathbf{D}}_r^H \Pi_{\tilde{\mathbf{A}}}^\perp \tilde{\mathbf{D}}_r) \odot (\mathbf{R}_s \tilde{\mathbf{A}}^H \tilde{\mathbf{R}}^{-1} \tilde{\mathbf{A}} \mathbf{R}_s)^T\right\} \quad (\text{A49})$$

$$\mathbf{H}_\theta = 2\text{Re}\left\{(\tilde{\mathbf{D}}_\theta^H \Pi_{\tilde{\mathbf{A}}}^\perp) \odot (\tilde{\mathbf{R}}^{-1} \tilde{\mathbf{A}} \mathbf{R}_s)^T\right\} \quad (\text{A50})$$

$$\mathbf{H}_r = 2\text{Re}\left\{(\tilde{\mathbf{D}}_r^H \Pi_{\tilde{\mathbf{A}}}^\perp) \odot (\tilde{\mathbf{R}}^{-1} \tilde{\mathbf{A}} \mathbf{R}_s)^T\right\} \quad (\text{A51})$$

$$\mathbf{K} = 2\text{Re}\left\{\Pi_{\tilde{\mathbf{A}}}^\perp \odot \tilde{\mathbf{R}}^{-T}\right\} - \Pi_{\tilde{\mathbf{A}}}^\perp \odot (\Pi_{\tilde{\mathbf{A}}}^\perp)^T \quad (\text{A52})$$

where $\tilde{\mathbf{D}} = [\tilde{\mathbf{D}}_\theta, \tilde{\mathbf{D}}_r]$, $\tilde{\mathbf{D}}_\theta = \mathbf{Q}^{-1/2} \mathbf{D}_\theta$, $\tilde{\mathbf{D}}_r = \mathbf{Q}^{-1/2} \mathbf{D}_r$, $\mathbf{D}_\theta \triangleq [\mathbf{d}_{\theta_1}, \mathbf{d}_{\theta_2}, \dots, \mathbf{d}_{\theta_K}]$, and $\mathbf{D}_r \triangleq [\mathbf{d}_{r_1}, \mathbf{d}_{r_2}, \dots, \mathbf{d}_{r_K}]$.

Finally, by substituting (A46)–(A52) into (A15), we can establish CRB(ζ) in (39) immediately. ■

ACKNOWLEDGMENT

The authors would like to thank the anonymous reviewers and the associate editor Prof. A.-K. Seghouane for their helpful criticism, comments, and suggestions that improved the manuscript.

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